Synchrotron X-ray structures in supernova remnants: particle acceleration probe

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Some R&D in the fields founded by V.L.Ginzburg and S.I.Syrovatskii Particle acceleration and the CR origin

Strong magnetic field amplification in cosmic particle accelerators and synchrotron X-ray images of SNRs

Non-linear mechanisms of efficient conversion of SN kinetic energy to relativistic components



Galactic cosmic-rays and SNR's

The power law, up to the "knee" at 10¹⁵ eV, is explained by diffusive shock acceleration at supernovae blast waves

Lagage and Cesarsky estimated the maximum energy to be less than 10¹⁴ eV assuming Bohm diffusion in a parallel shock geometry without magnetic field amplification

A higher maximum energy is expected for a *quasi-perpendicular* shock



X-ray synchrotron & B-field amplification



from Jacco Vink 2009

Chandra profiles by A.Bamba



measurements of the width of synchrotron X-ray filaments ~ 0.01pc STRONG MAGNETIC FIELD AMPLIFICATION > 20 μ G Electron energies >> 1 TeV

Particle accelerators – LHC– SN1006









- SN 1572
- SN type: la
- distance: ~3 kpc
- radius: ~3.7 pc

★ Cassiopeia A

- SN ~1680
- SN type: IIb
- distance: ~3.4 kpc
- radius: ~2.5 pc

X-ray Images (Chandra)

Most parameters are reasonably well known. → largely help us interpret gamma-ray results.

Uchiyamaa Fermi Symp 20111



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Warren+05 Chandra shock heated ejecta

X-ray/radio radial profile



B₂ = 0.1-0.2 mG is inferred from the width of X-ray filaments

synchrotron X-rays

Chandra 4-6 keV Image of Tycho's SNR



Eriksen + 2011

Tycho: New GeV Detection



Gamma-ray Space Telescope



See a poster by Fermi-LAT Collaborati	on
(Naumann-Godo+)	



Photon index = 2.3 ± 0.1 (favors hadronic origin)

6-8% of E _{SN}	
transferred to	CRs.

Case	D _{kpc}	n _H [cm ⁻³]	E _{sn} [10 ⁵¹ erg]	E _{p,tot} [10 ⁵¹ erg]	K _{ep}
Far	3.50	0.24	2.0	0.150	4.5x10 ⁻⁴
Nearby	2.78	0.30	1.0	0.061	7.0x10 ⁻⁴

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Tycho: Recent TeV Detection

Gamma-ray Space Telescope



Flux(>1 TeV) ~ 1% Crab 5.0σ detection (post-trial)

B-field constraint put by X-ray does *not* contradict IC origin.

Fermi-LAT can test "leptonic vs hadronic"



Ultra-relativistic Particle Acceleration in collisionless shocks by Fermi mechanism







Particles make nearly elastic collisions with background plasma
→ gain energy when cross shock → bulk kinetic energy of converging flows put into individual particle energy



In efficient acceleration, <u>entire particle spectrum</u> must be described consistently, including escaping particles → much harder mathematically BUT, connects photon emission across spectrum from radio to γ-rays Nonlinear Models of Ultrarelativistic Particle Acceleration with CR driven instabilities



- Our model simulates particle acceleration, turbulence generation and shocked flow all consistently with each other;
- A Monte Carlo (MC) code describes particle transport and acceleration;
- Diffusion coefficient used in the MC code coupled to turbulence spectrum;
- Turbulence generation driven according to particle transport simulated in MC.

Method

The Nonlinear Model

- Particle transport modeled with a Monte Carlo simulation;
- Analytic, semi-phenomenological description for magnetic field amplification, self-consistently coupled to CR distribution and MHD flow;
- Fundamental conservation laws used to iteratively derive a nonlinear shock modification that conserves mass, momentum and energy;

Reasoning

- We describe a large dynamic range in turbulence scales and particle energies;
- Elements of the model tested against spacecraft observations of heliospheric shocks;
- Works for highly anisotropic particle distributions (particle escape and injection; large gradients of *u* and *B*).
- Ability to incorporate non-diffusive particle transport (future work).

Evolution of Waves in the Precursor

Definitions

We describe turbulence by W(x, k) – spectral energy density of turbulent fluctuations, and separate it into

$$W = W_M + W_K = \sum_{i \in \text{modes}} W_M^{(i)} + \sum_{i \in \text{modes}} W_K^{(i)}.$$

 W_M – magnetic fluctuations, W_K – associated plasma velocity fluctuations, and (*i*) runs over the three types of waves (A – Alfvén waves, B – Bell's modes, C – s modes).

Equations

Evolution for each mode is given by the equation for $W^{(i)} = W_M^{(i)} + W_K^{(i)}$:

$$u\frac{\partial W^{(i)}}{\partial x} = \gamma^{(i)}W^{(i)} - L^{(i)} + \left[-\alpha^{(i)}W^{(i)} + \frac{\partial}{\partial k}\left(kW^{(i)}\right)\right]\frac{du}{dx} - \frac{\partial\Pi^{(i)}}{\partial k}$$
(1)

CR driven modes



Kulsrud 69, Bell' 04, AB+05, Marcowith+ 06, Zirakashvili & Ptuskin 08, AB+11

CR modified shock



• Vladimirov, Bykov & Ellison, 2009. ApJ, v. 703, L29

Particle Spectra





[•] Vladimirov, Bykov & Ellison, 2009. ApJ, v. 703, L29

Magnetic Turbulence in the Geospace Environment



Fig. 11 Magnetic power distribution in k space, as a function of wavevector components (in units of the proton gyroradius) for directions along the background magnetic field (*left*), and along the plasma velocity (nearly perpendicular to the magnetic field) (adapted from Sahraoui et al. 2006). Reprinted with permission from Sahraoui et al. (2006). Copyright 2006 by the American Physical Society

Zimbardo + 2010

No parallel cascade...

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e.g. Ginzburg and Syrovatskii 1969



FIG. 5. Oscillation ellipse of the electric vector in a wave radiated by particles moving in a magnetic field, where the charge is taken as a positive. For negatively charged particles (electrons) the direction of rotation is opposite to that shown. The plane K is the plane of the figure (the plane perpendicular to the direction of the radiation or, equivalently, to the direction of the observer), and l_1 and l_2 are two mutually orthogonal unit vectors in the plane of the figure, of which l_2 is directed along the projection of the magnetic field H on the plane K.

Synchrotron Radiation Stockes Parameters:

$$\hat{\tilde{S}} = \begin{pmatrix} \tilde{I}(\mathbf{r}, t, \nu) \\ \tilde{Q}(\mathbf{r}, t, \nu) \\ \tilde{U}(\mathbf{r}, t, \nu) \\ \tilde{V}(\mathbf{r}, t, \nu) \end{pmatrix} = \begin{pmatrix} p_{\nu}^{(1)} + p_{\nu}^{(2)} \\ (p_{\nu}^{(1)} - p_{\nu}^{(2)}) \cdot \cos 2\chi \\ (p_{\nu}^{(1)} - p_{\nu}^{(2)}) \cdot \sin 2\chi \\ (p_{\nu}^{(1)} - p_{\nu}^{(2)}) \cdot \sin 2\beta \end{pmatrix}$$

$$\hat{S}(\mathbf{R}_{\perp},t,\nu) = \int dl \, d\gamma \, N(\mathbf{r},\gamma,t') \, \hat{\tilde{S}}(\mathbf{r},t',\nu,\gamma), \quad t' = t - |\mathbf{r} - \mathbf{R}_{\perp}|/c.$$

To construct the synchrotron emission image we simulated stochastic magnetic field in a SNR shell

х



Time evolution. Lightcurves.





Synchrotron X-ray images at energies 0.5, 5, 20, 50 keV (from left to right). Dot like feature D1 is clearly seen at high energies and it is smeared in at low energies. Left panels show lightcurves of D1 feature at 5 keV (solid curve), 20 keV (dashed curve) and 50 keV (dot-dashed curve).

AB+ ApJL 2008



Uchiyama Aharonian et al. 2007

Nonthermal clump "lifetime" ~ 1yr It can be produced with magnetic field well below 1 mG ..

X-ray Polarization at 5 keV



X-ray Polarization at 50 keV



 $\delta = 1.0$

X-ray strips in Tycho's SNR (Eriksen etal 2011)



FIG. 1.— Chandra X-ray 4.0–6.0 keV image of the Tycho supernova remnant, smoothed with a $\sim 0.75''$ Gaussian and displayed with an *arcsinh* scaling, showing various regions of striping in the nonthermal emission. Clockwise from the upper right: a) The main western stripes discussed in this Letter; b) A fainter ensemble of stripes; c) a previously-known bright arc of non-thermal emission, with our newly discovered streamers; d) filaments of "rippled sheet" morphology common in optical observations of middle-aged SNRs.

Chandra 4-6 keV X-rays



Magnetic Fluctuation Spectra



Chandra 4-6 keV Image of Tycho's SNR



Eriksen + 2011



Polarization fraction

AB+ ApJL v735, L40, 2011

No simple explanation of strips !

➔ Many shock and turbulence properties must come together to produce coherent structure on this scale.

Strong predictions: Quasi-perpendicular upstream B-field

Strong linear polarization in strips