



Quest for the QCD phase diagram in extreme environments



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How “extreme” typically?



High Temperature

up to $T \sim \Lambda_{\text{QCD}} \sim 200 \text{ MeV}$

Relativistic Heavy-Ion Collision

High Baryon Density

up to $\rho_B \sim (\Lambda_{\text{QCD}})^3 \sim 1 \text{ fm}^{-3}$

Relativistic Heavy-Ion Collision, Neutron Star

Strong Magnetic Field

up to $eB \sim (\Lambda_{\text{QCD}})^2 \sim 10^{18} \text{ gauss}$

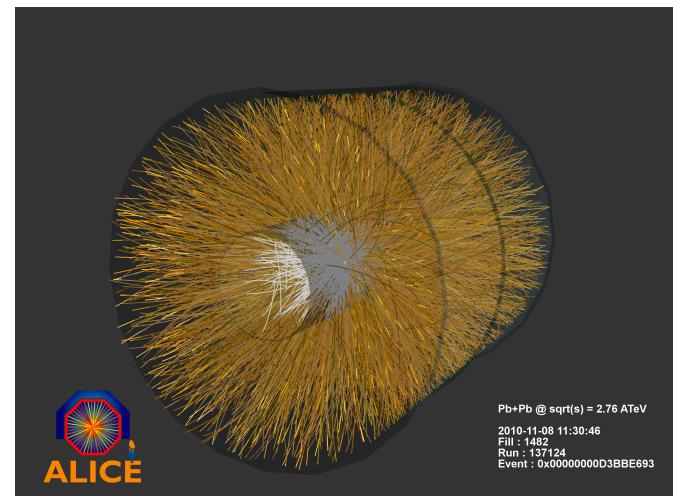
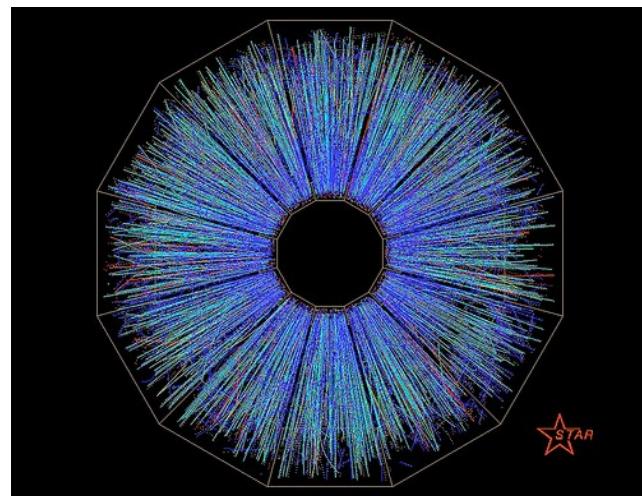
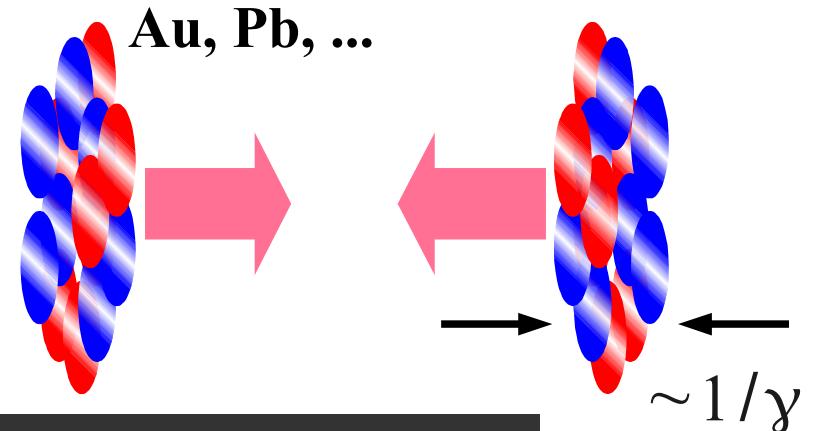
Relativistic Heavy-Ion Collision, Neutron Star

Relativistic Heavy-Ion Collision



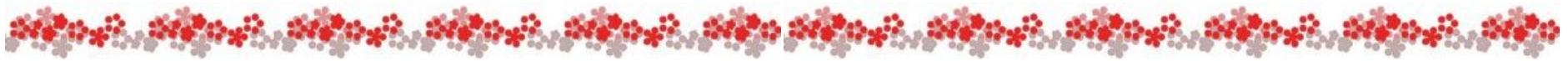
LHC: $\sqrt{s_{NN}} = 2.7 \text{ TeV} \rightarrow \gamma \sim 1400$

RHIC: $\sqrt{s_{NN}} = 200 \text{ GeV} \rightarrow \gamma \sim 100$

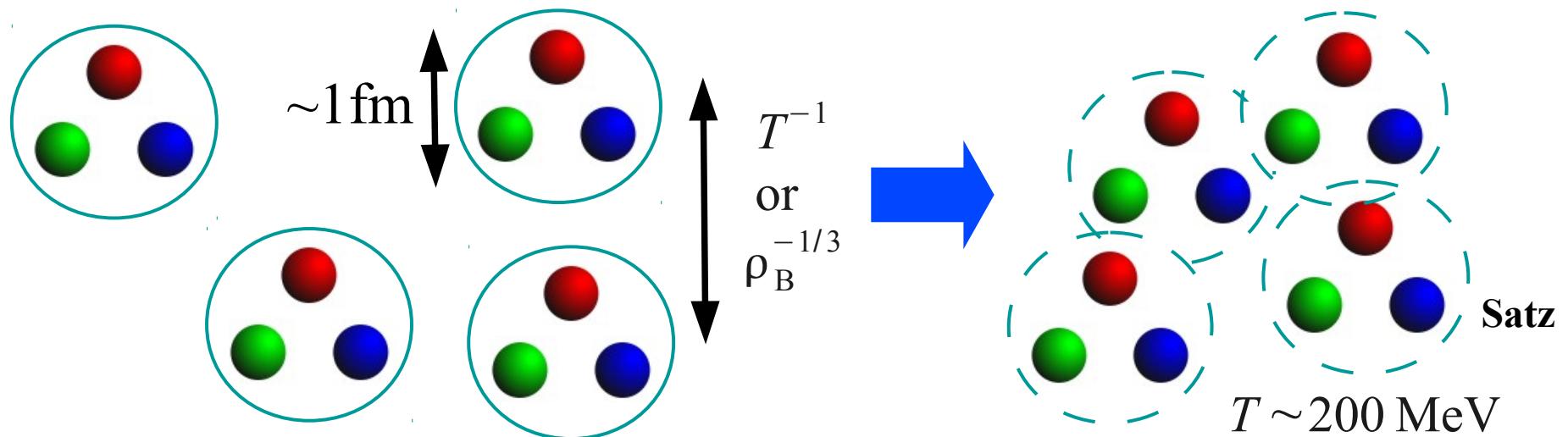


Thermalization achieved (elliptic flow by a hydro-model)
Initial temperature $\sim 400 \text{ MeV}$ (distribution of direct photon)

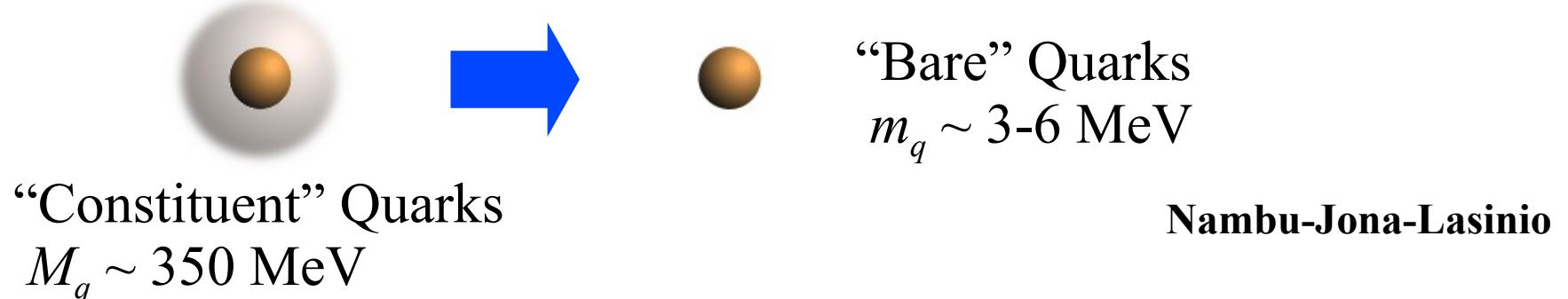
Two Major Phase Transitions in QCD



Quark Deconfinement Transition (Polyakov Loop)

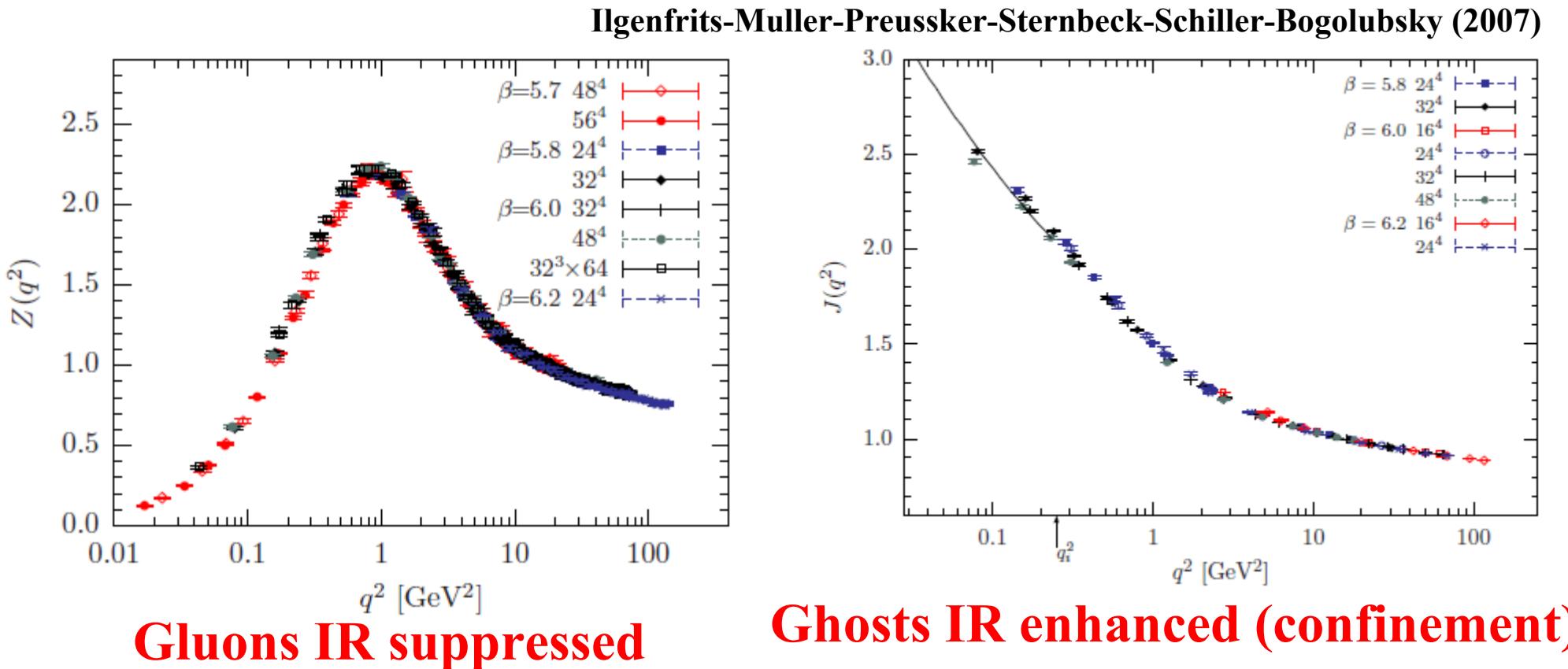


Chiral Phase Transition (Chiral Condensate)

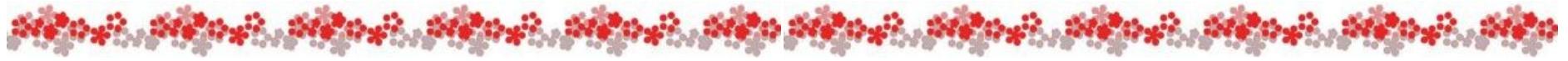


How to understand deconfinement?

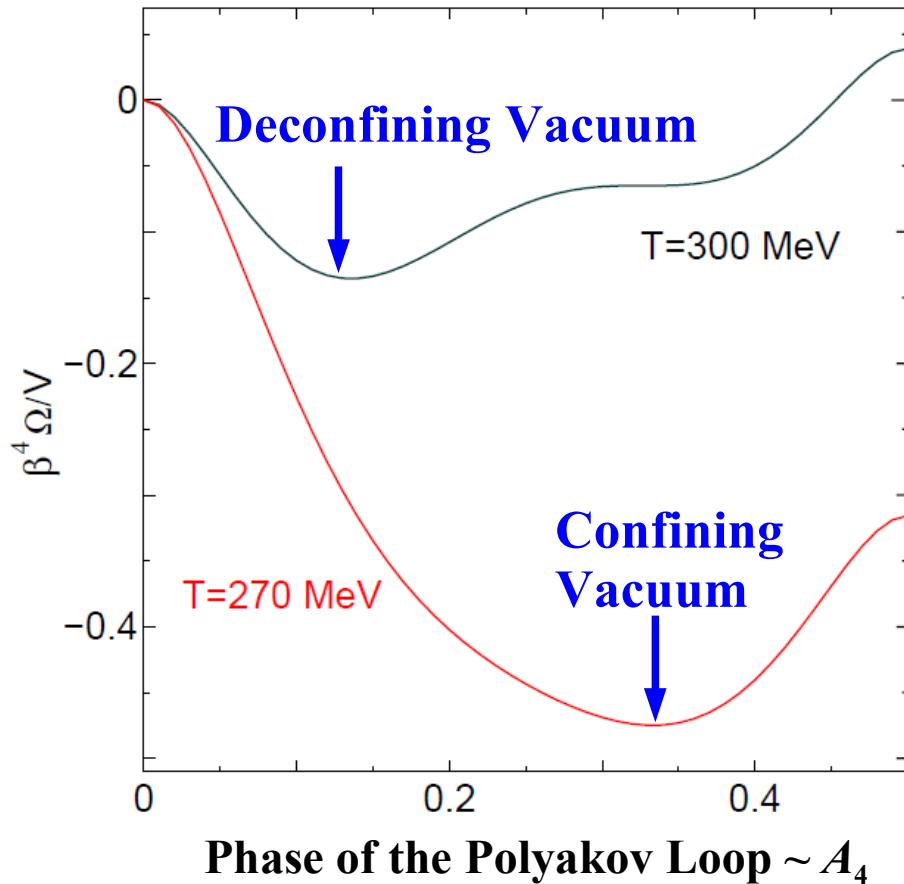
Confinement understood from the non-perturbative propagators of gluons and ghosts in the Landau gauge



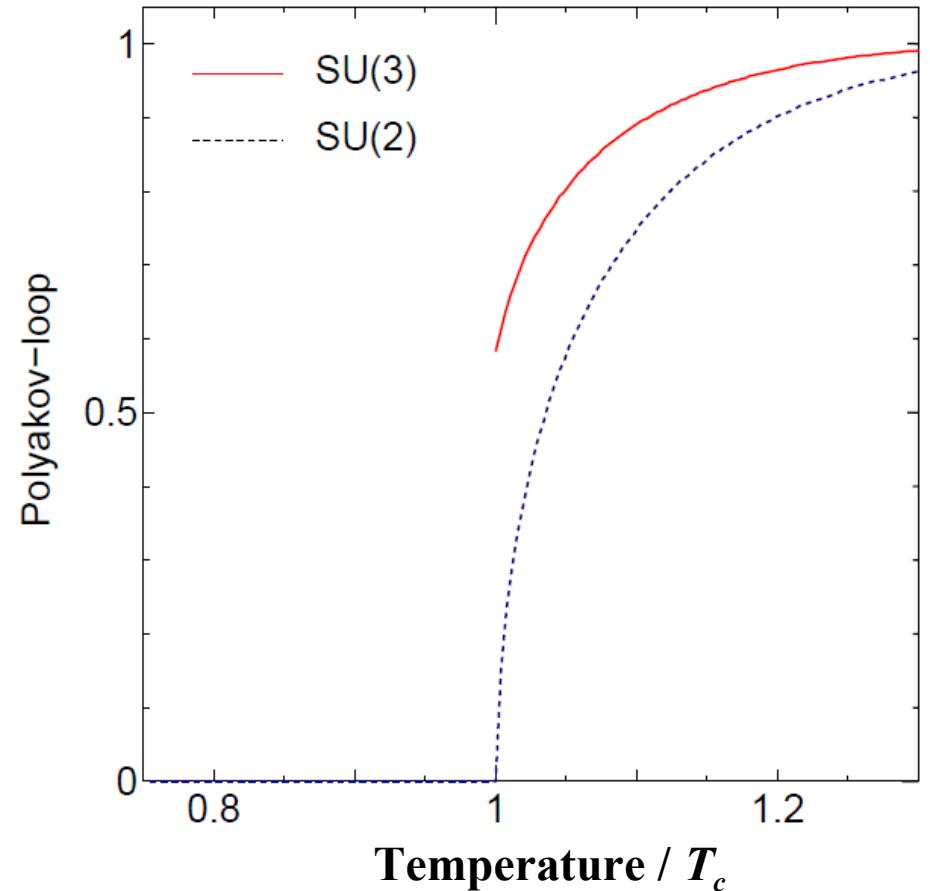
Deconfinement from Propagators



$$\ln Z = -\frac{1}{2} \text{tr} \ln D_A^{-1}(A_4) + \text{tr} \ln D_C^{-1}(A_4) + \dots$$

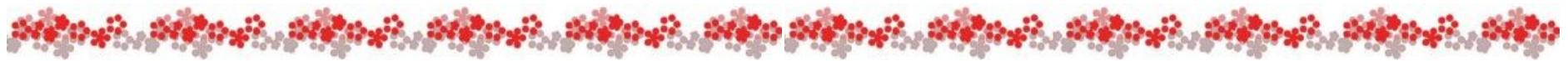


Confinement \rightarrow Deconfinement



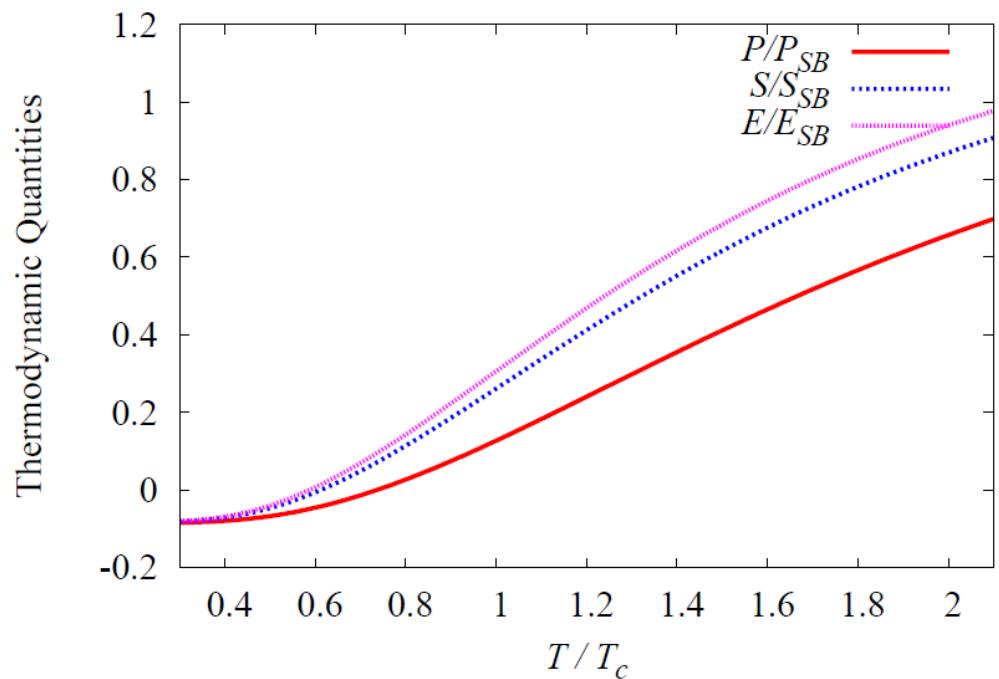
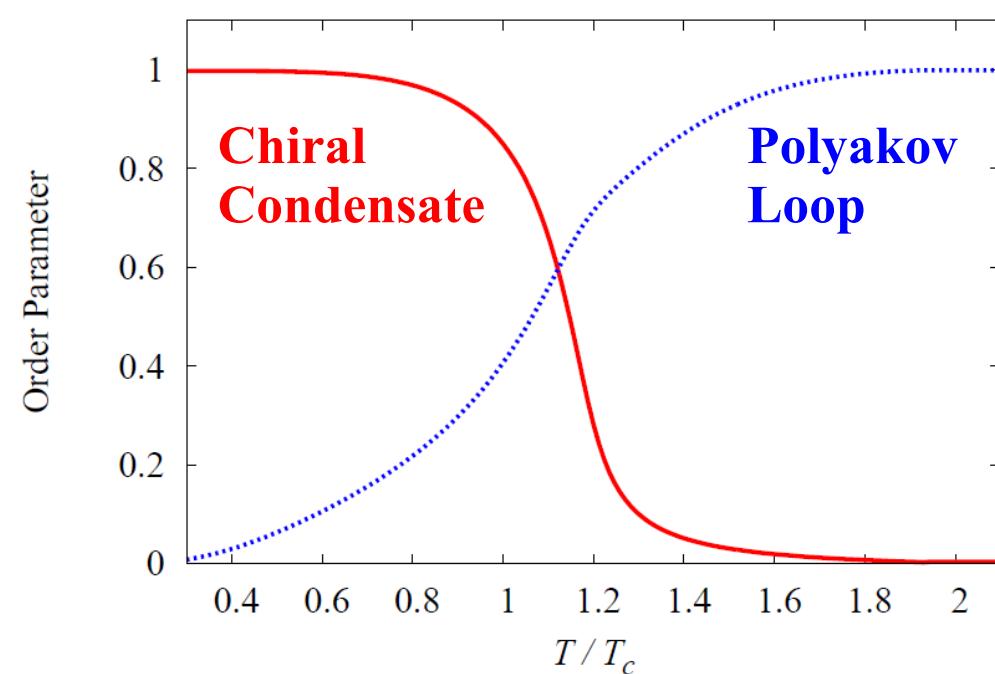
Braun-Gies-Pawlowski, KF-Kashiwa

Coupled with Dynamical Quarks



$$\ln Z_F = \text{tr} \ln (1 + e^{-(E-\mu)/T}) + \text{tr} \ln (1 + e^{-(E+\mu)/T})$$

➡ $\text{tr} \ln (1 + L e^{-(E-\mu)/T}) + \text{tr} \ln (1 + L^\dagger e^{-(E+\mu)/T})$



Polyakov-loop potential determined unambiguously

KF-Kashiwa (2012)

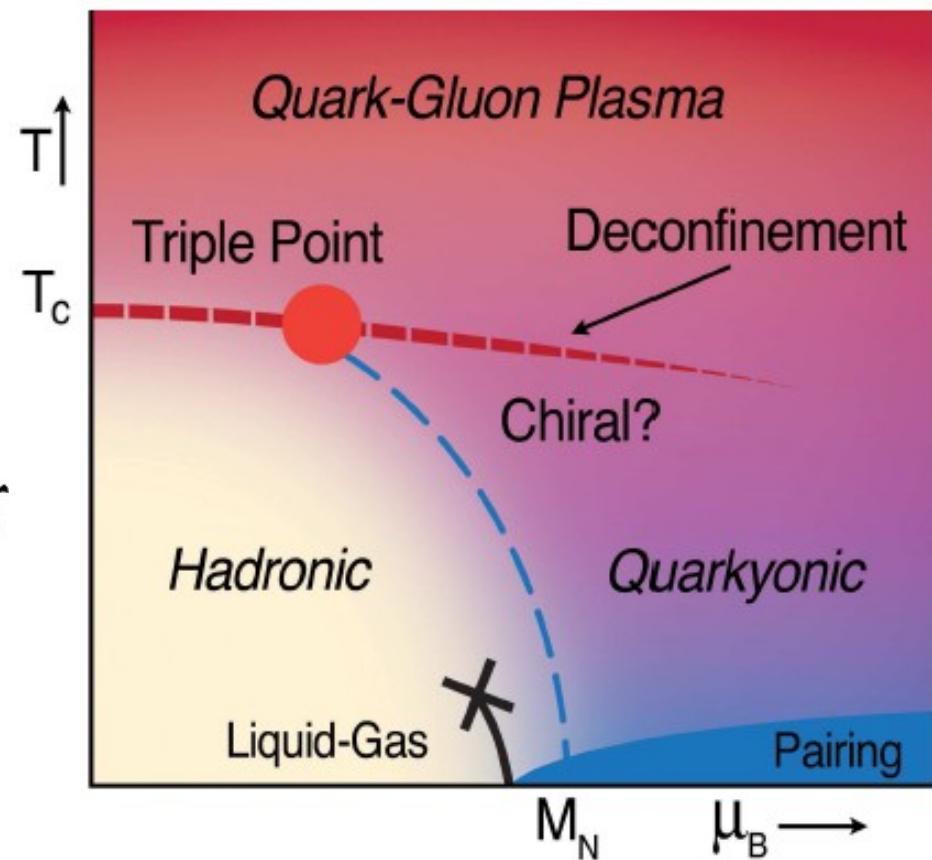
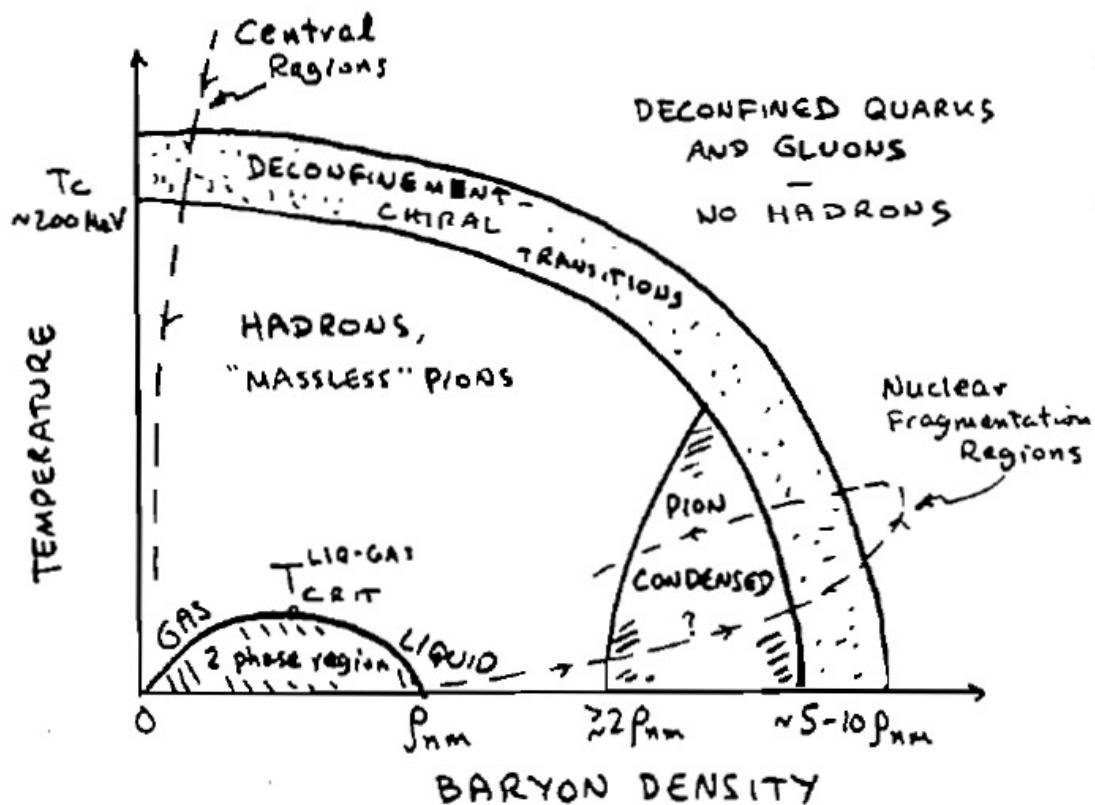
Phase Diagrams including Density

Prototype in 1983

and

Update in 2009

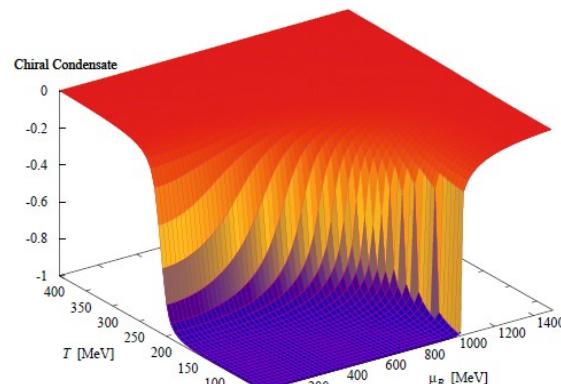
PHASE DIAGRAM OF NUCLEAR MATTER.



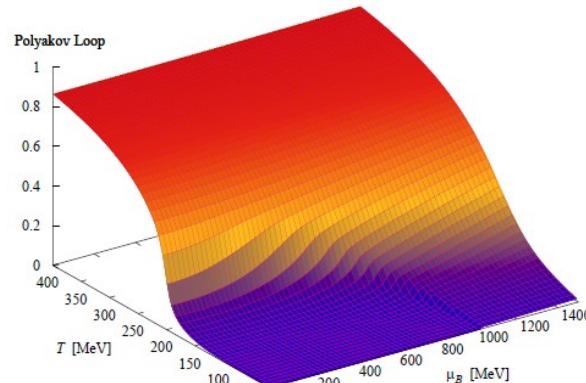
Modern View

Effective Model Results

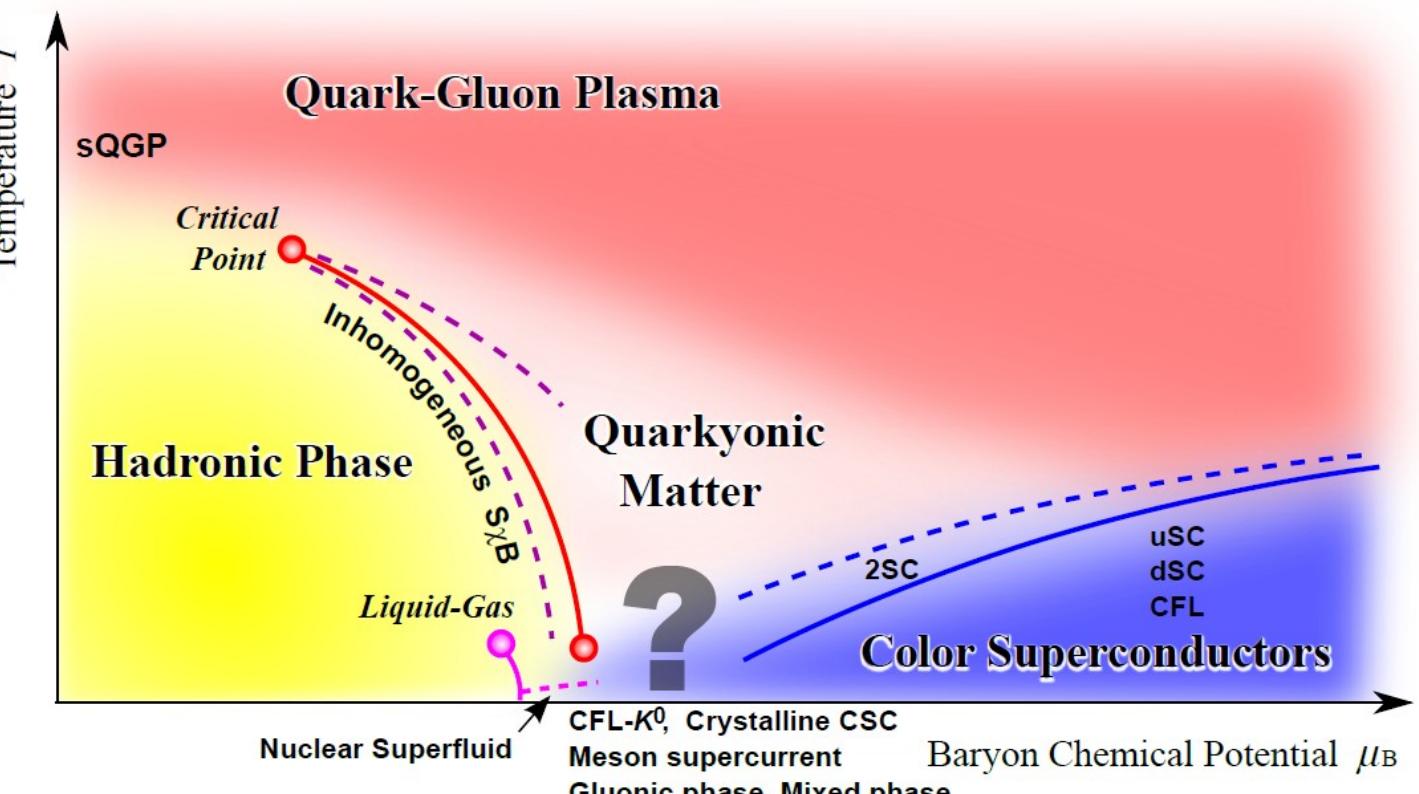
Chiral Condensate



Polyakov Loop



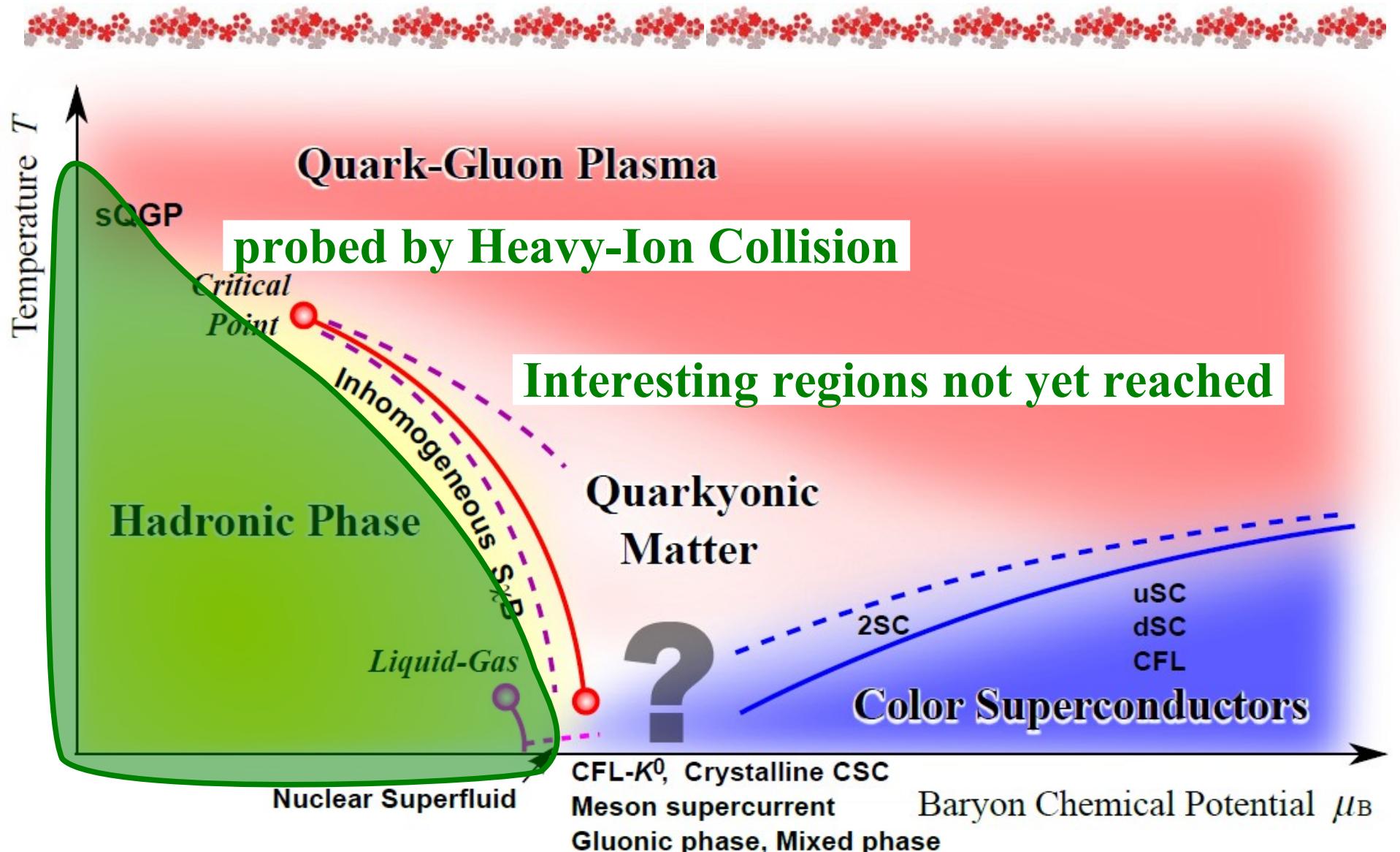
Conjectured Phase Structure



KF (2008)

KF-Hatsuda (2010)

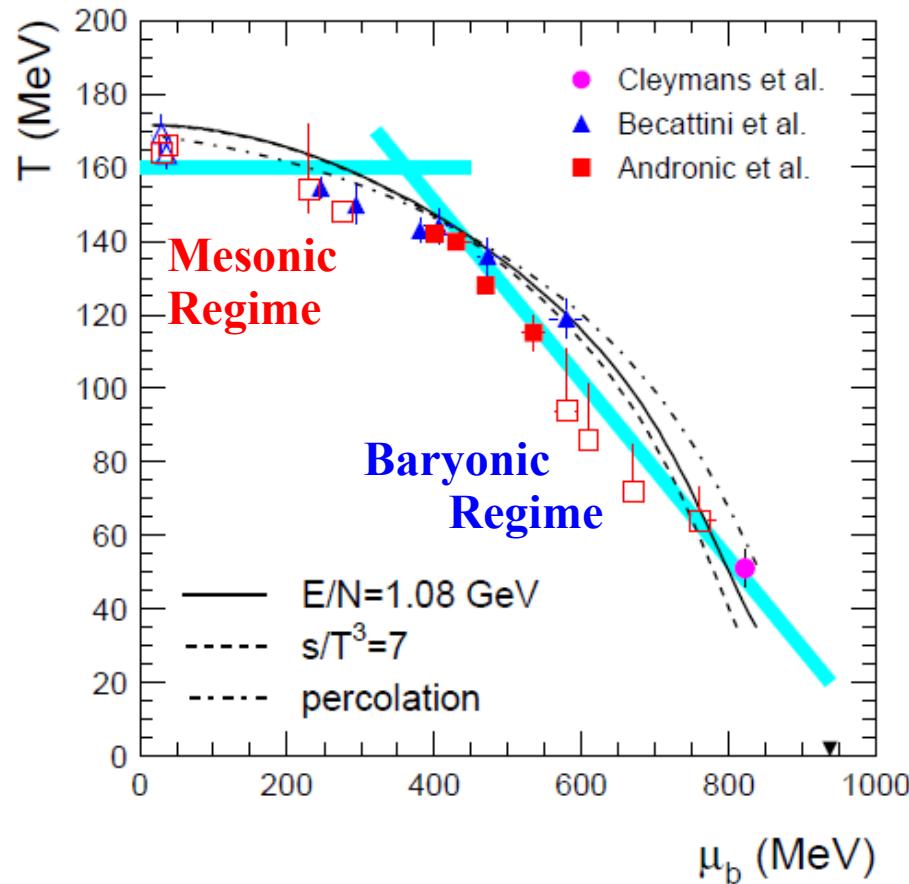
Experimentally Confirmed



KF-Hatsuda (2010)

Experimental Data

Freeze-out points are located by the particle yields
 Two regimes in **meson-dominance** and **baryon-dominance**



Mesonic Hagedorn Transition

$$Z \sim \int dm \rho(m) e^{-m/T}$$

$$\rho(m) \sim e^{m/T_H}$$

$$T_c = T_H$$

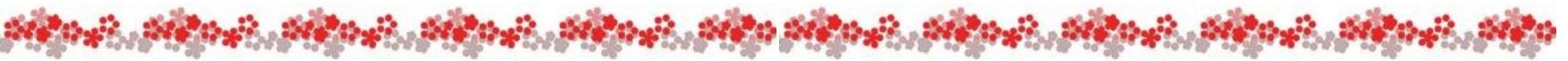
Baryonic Hagedorn Transition

$$Z \sim \int dm \rho_B(m) e^{-(m_B - \mu_B)/T}$$

$$\rho(m) \sim e^{m_B/T_B}$$

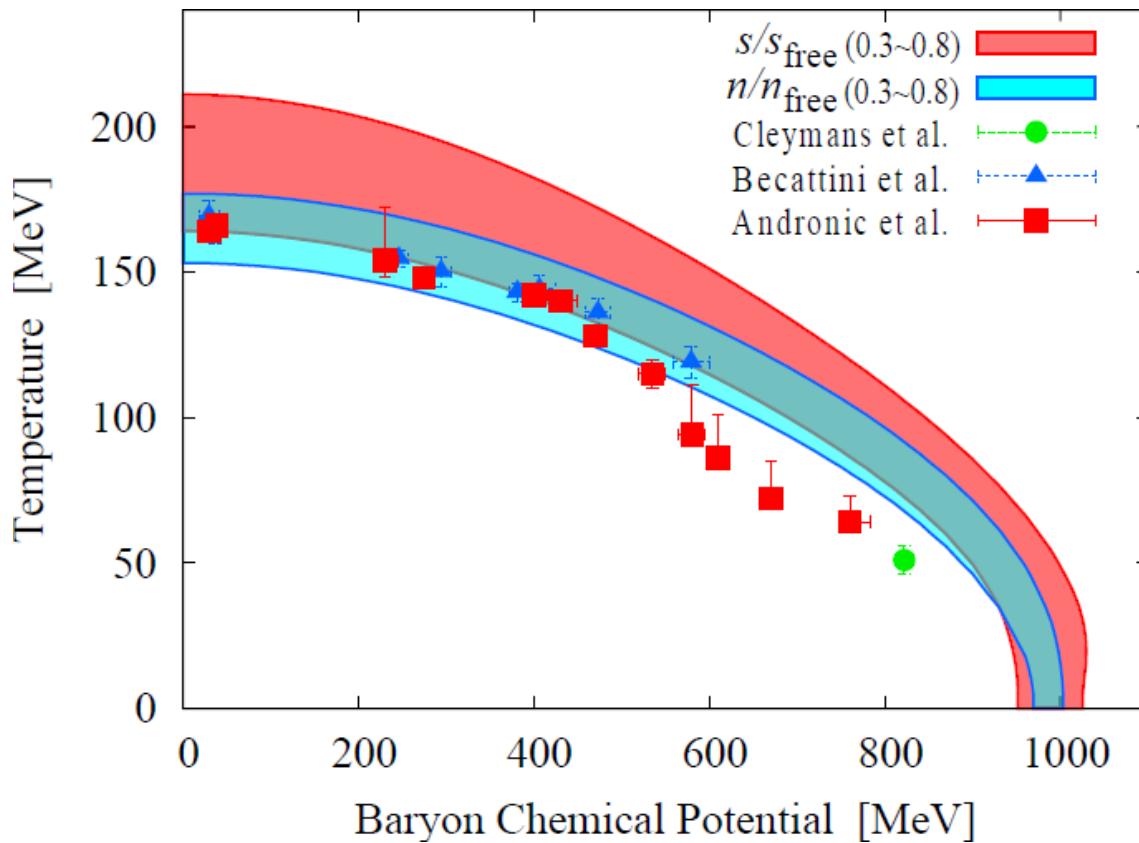
$$T_c = (1 - \mu_B/m_B) T_B$$

Thermodynamics



Statistical Model Interpretation

KF (2010)

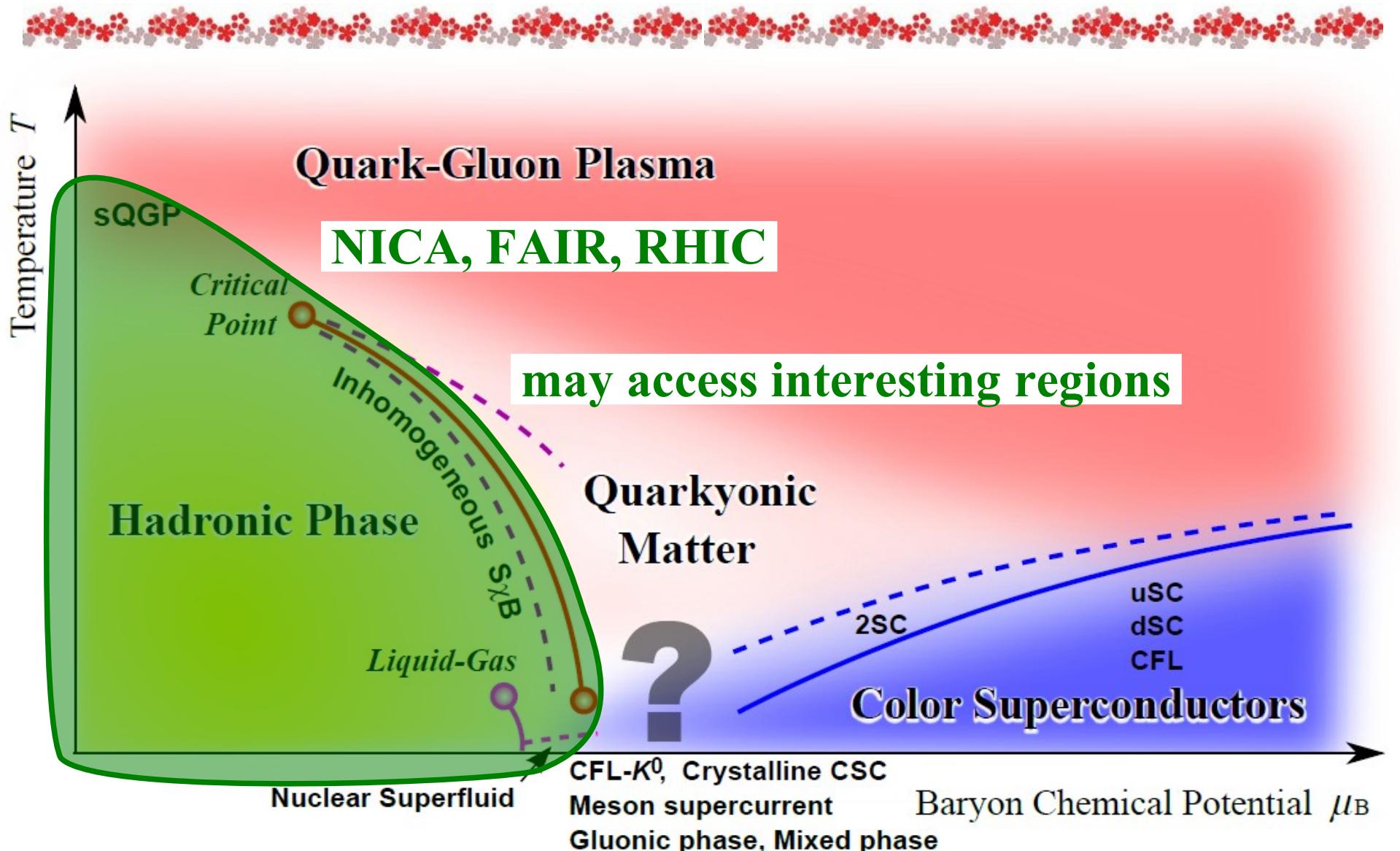


Gluon Deconfinement
~ Increasing entropy

Quark Deconfinement
~ Increasing density

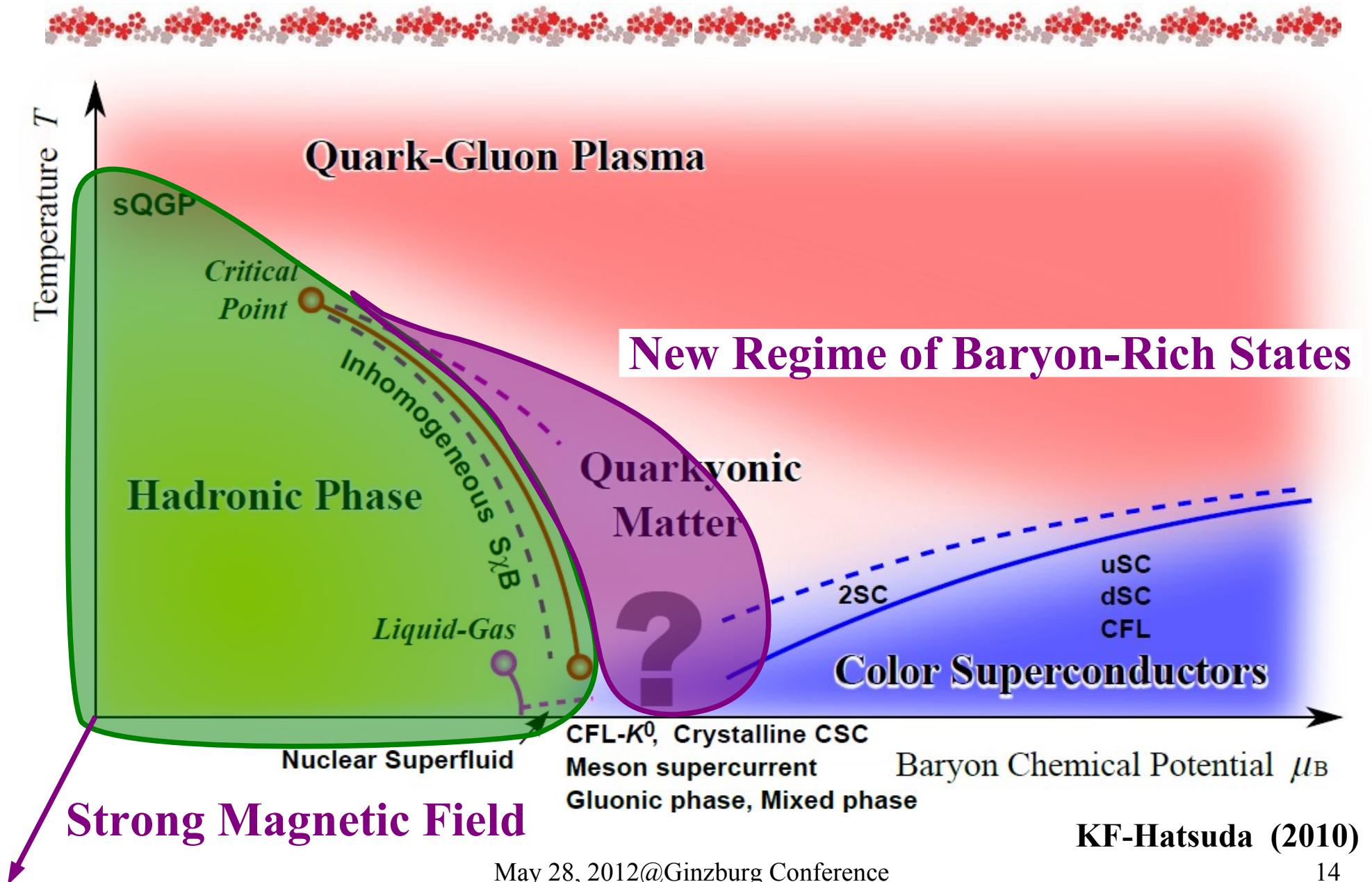
Thermodynamic quantities
taken over by (quasi-)gluons
and (quasi-)quarks
(beyond the Hagedorn limit)

Experimentally Expected

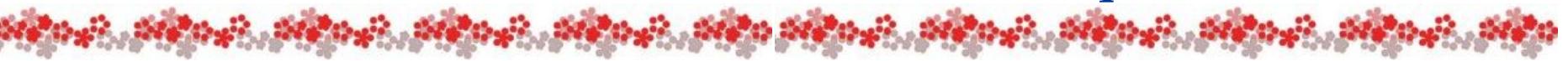


KF-Hatsuda (2010)

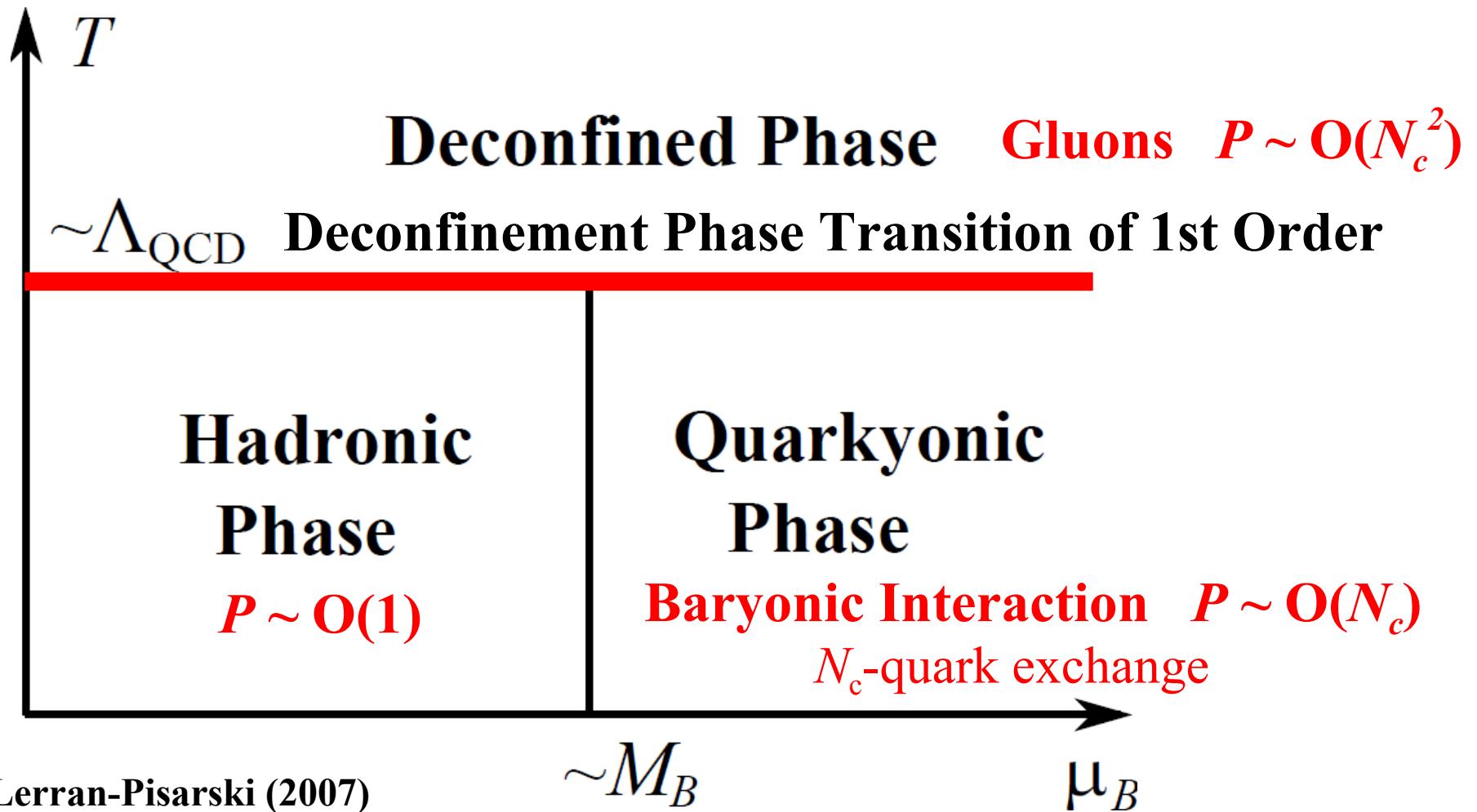
Theoretically Speculated



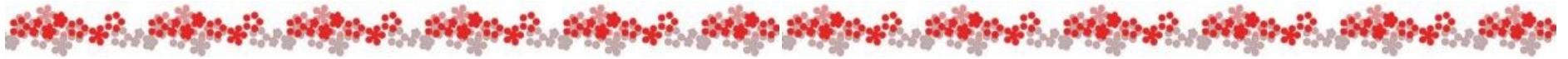
New Regime at Large μ_q and N_c



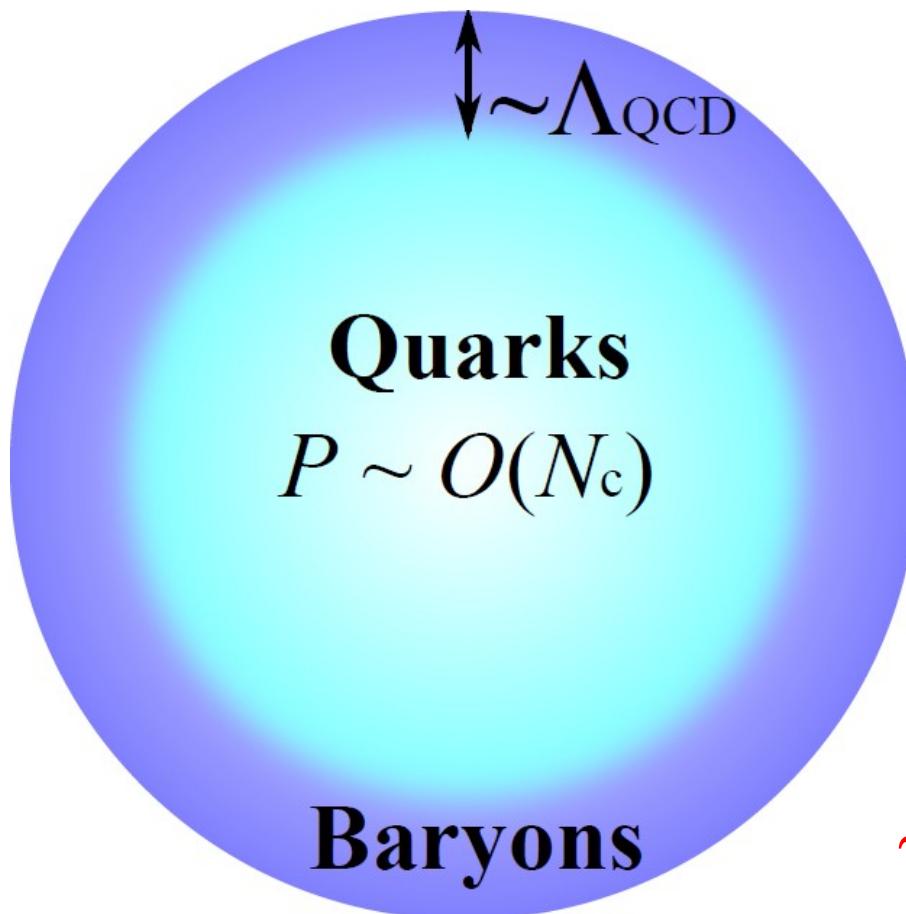
Phase Diagram of Large- N_c QCD



Quarkyonic Matter



Structure of the Fermi Sphere



**Ground state of
large- N_c quark matter
at $\mu_q \gg \Lambda_{\text{QCD}}$**

McLerran, Pisarski
Hidaka, Kojo

**Interacting Baryon Crystal
~ Quasi-quark Gas**

Quarkyonic Chiral Spiral ($\mu_q \gg \Lambda_{QCD}$)



**Choose one direction z with $p_z \sim \mu_q$ ($p_x, p_y \sim \Lambda_{QCD}$)
(1+1)D system effectively**

$$\bar{\psi} (i \gamma^z \partial_z + \mu \gamma^0) \psi \\ = \bar{\psi}' (i \gamma^z \partial_z) \psi' \quad \quad \psi = e^{i \gamma^0 \gamma^z \mu z} \psi'$$

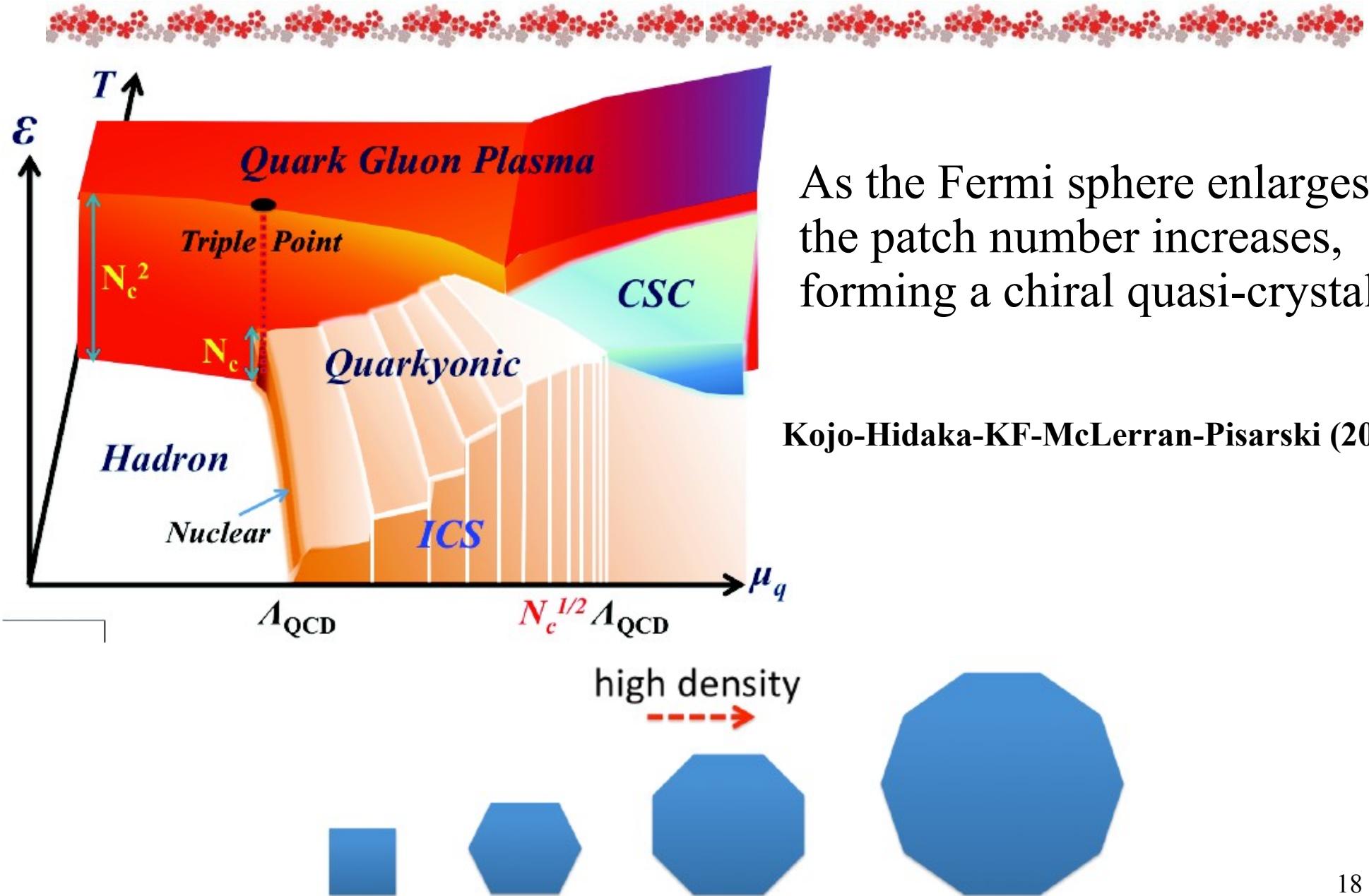
$\langle \bar{\psi}' \psi' \rangle$ = Homogeneous condensate at zero density

$$\langle \bar{\psi} \psi \rangle = \langle \bar{\psi}' \psi' \rangle \cos(2\mu z)$$

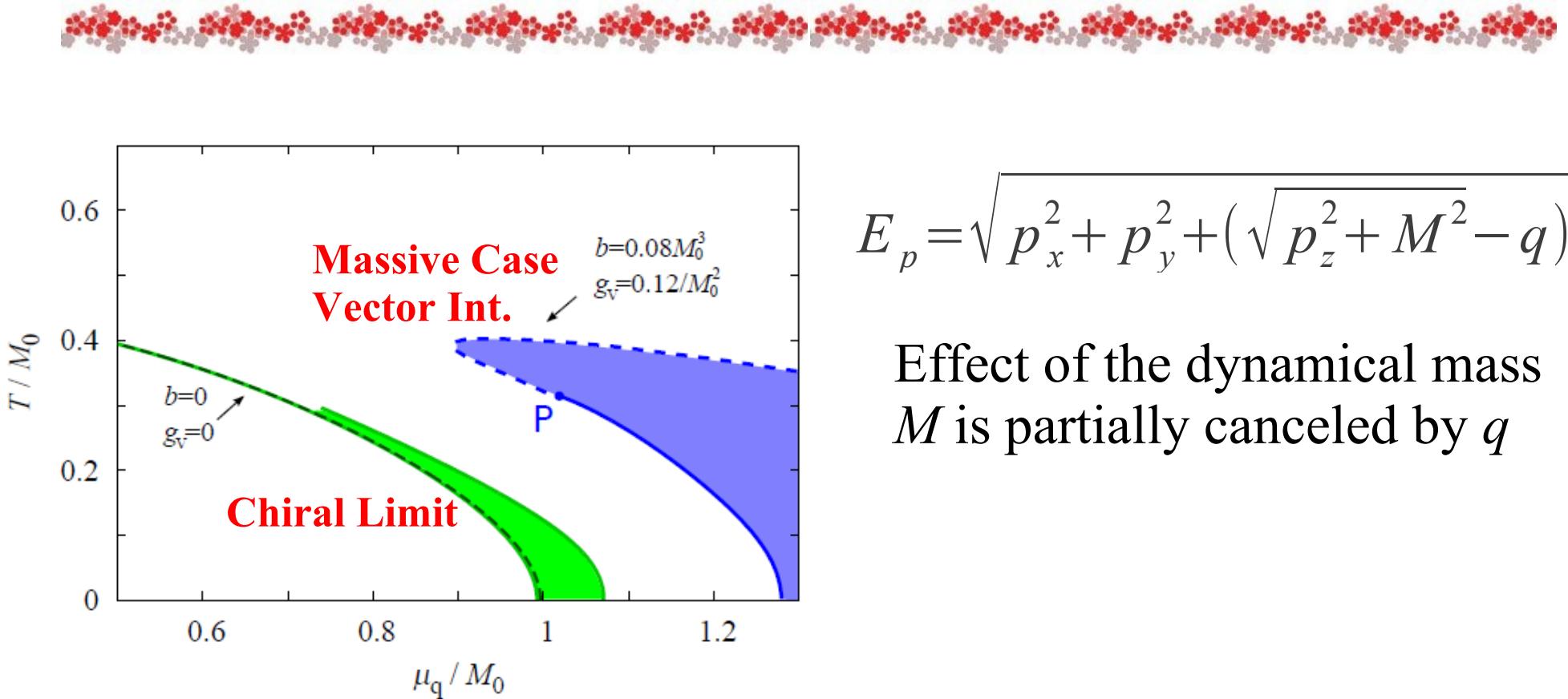
$$\langle \bar{\psi} \gamma^0 \gamma^z \psi \rangle = \langle \bar{\psi}' \psi' \rangle \sin(2\mu z)$$

This quasi-(1+1)D system forms “one patch”

Interweaving Chiral Spirals



Some Model Results



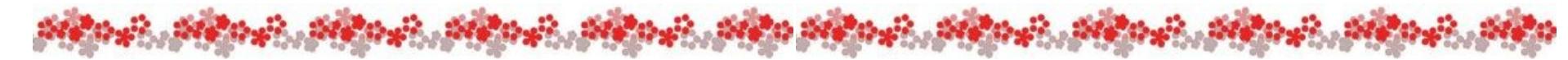
$$E_p = \sqrt{p_x^2 + p_y^2 + (\sqrt{p_z^2 + M^2} - q)^2}$$

Effect of the dynamical mass M is partially canceled by q

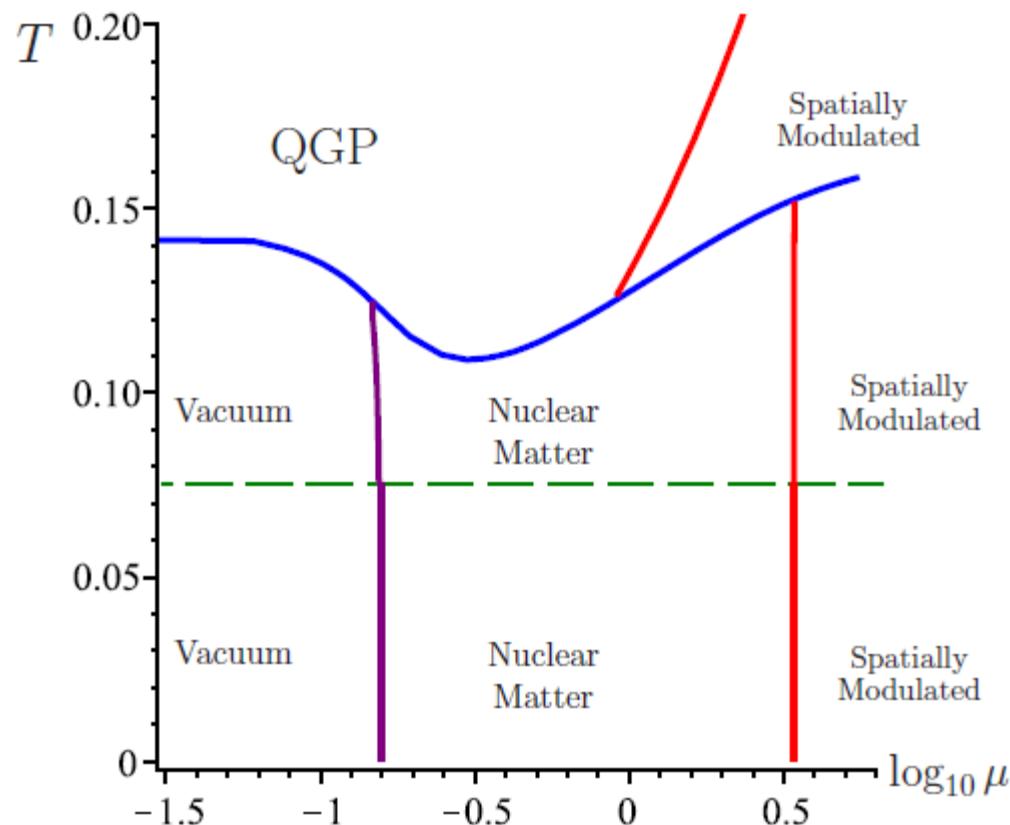
Even when N_c and μ_q are not infinitely large,
the chiral spiral is favored near the phase boundary
of chiral symmetry

Nakano-Tatsumi (2003), KF (2012)

Holographic Evidence



State-of-the-art phase diagram in holographic model



Instability to inhomogeneous states is seen (in a different way from QCD...)

Nakamura-Ooguri-Park, Chuang-Dai-Kawamoto-Lin-Yeh (2010)

Density Effect ~ Magnetic Field Effect



Energy dispersion relation in B

$$\omega^2 = p_z^2 + \underline{2|eB|(n+1/2)} + m^2 - 2s e B$$

Transverse motion = Harmonic Oscillator

Fermions ($s=1/2$) have zero mode – dominant at large B
Quasi-(1+1)D system is realized along the B direction.

Very strong B + Any μ_q → Chiral Spiral

Basar-Dunne-Kharzeev

Very strong B + Attractive Int.

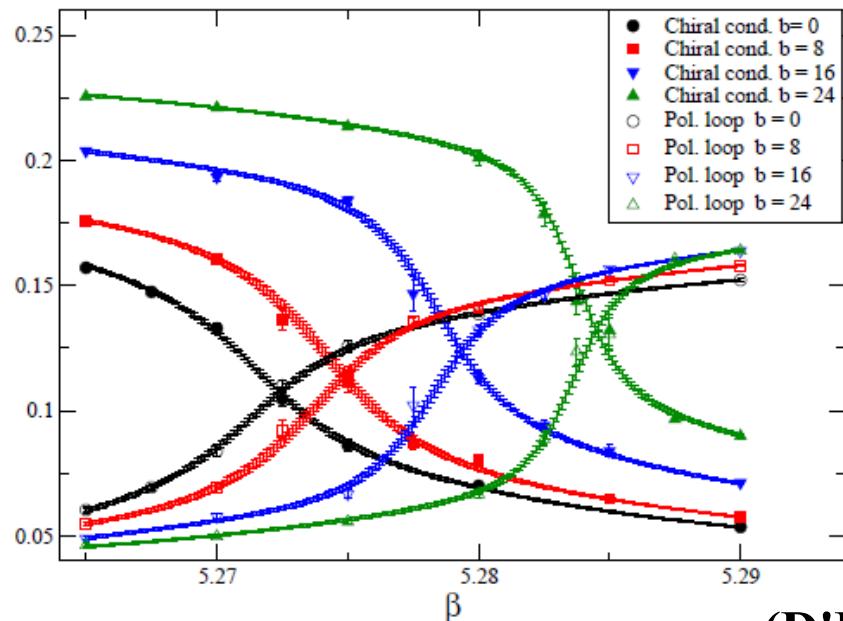
→ Cooper Instability → Magnetic Catalysis

Klimenko, Gyusynin-Miransky-Shovkovy

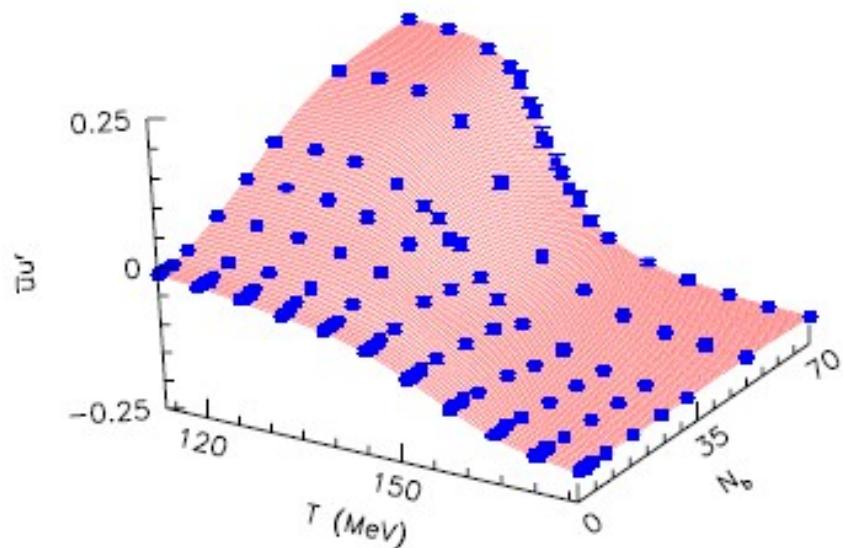
B Effect on the Phase Diagram



QCD phase transitions affected by B



(D'Elia et al)



(Fodor et al)

Monte-Carlo simulation is possible (no sign problem)

T_c increases or decreases?

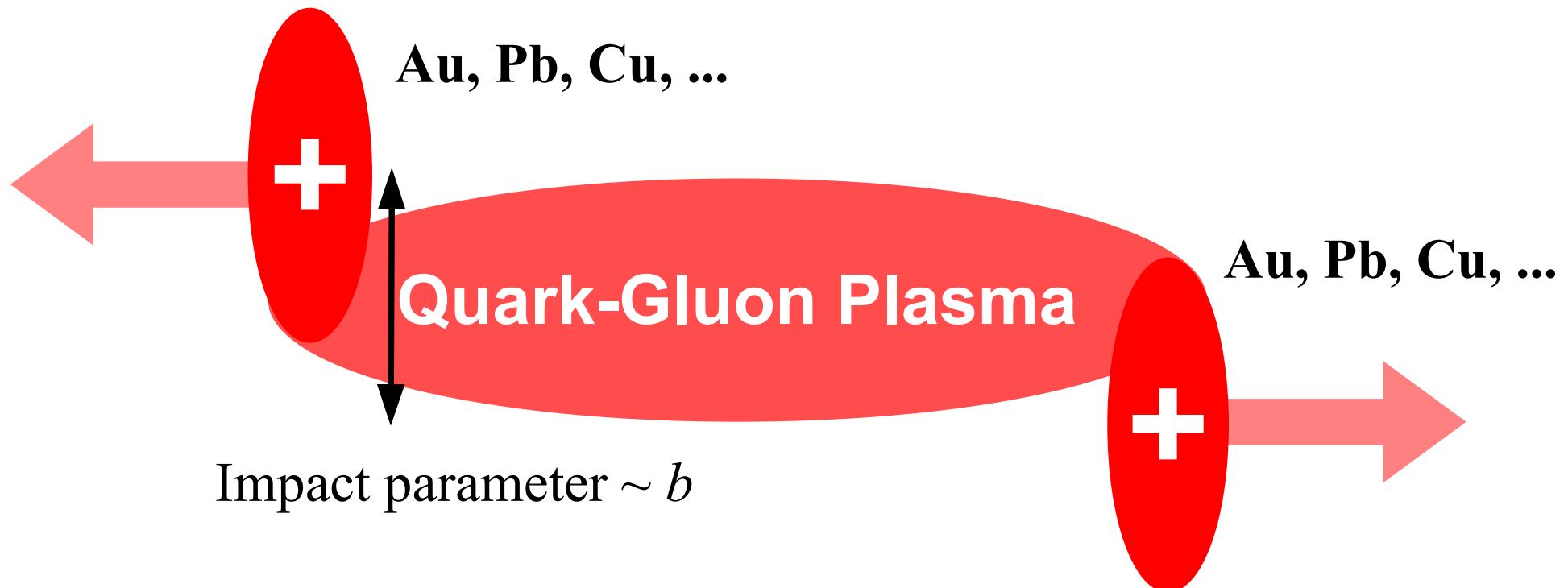
Contradictory results from two groups!

Not Only Theoretical Analogue

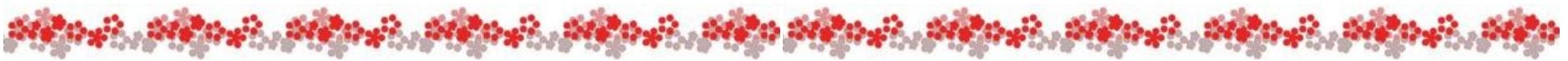


Relativistic Heavy-Ion Collision

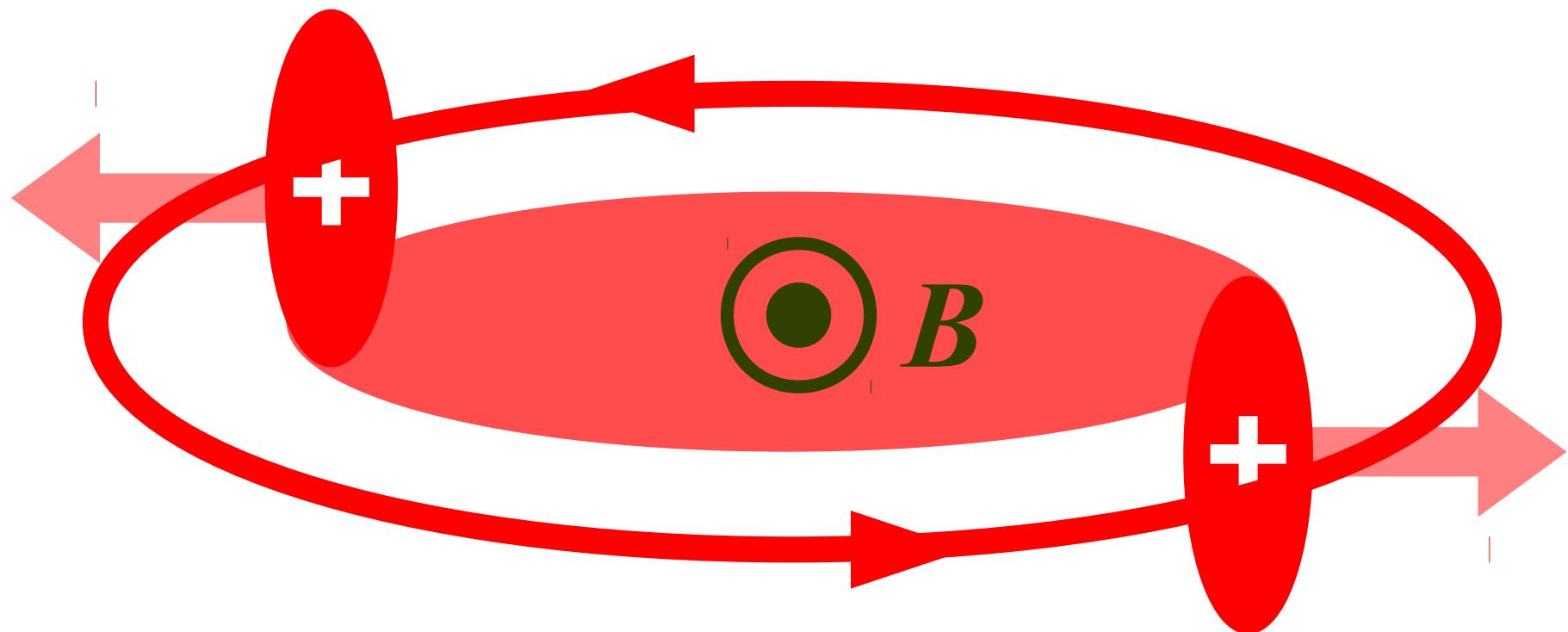
Moving almost at the speed of light



Origin of the Magnetic Field



Strong B generated due to Electrodynamics

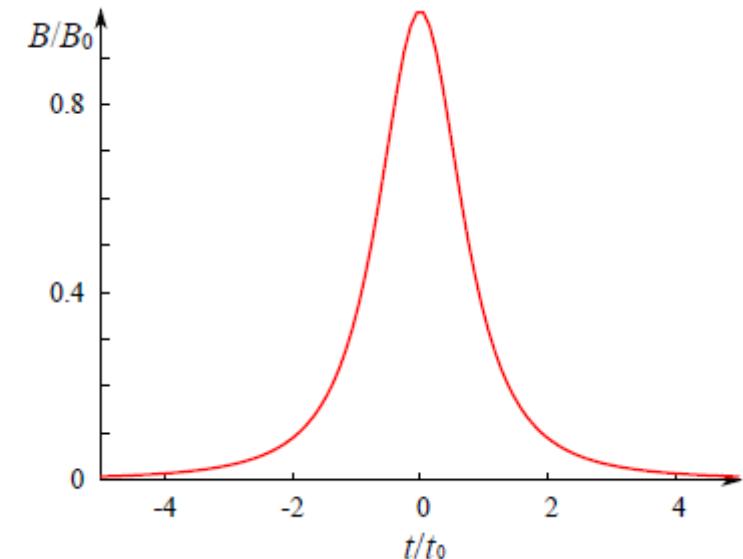
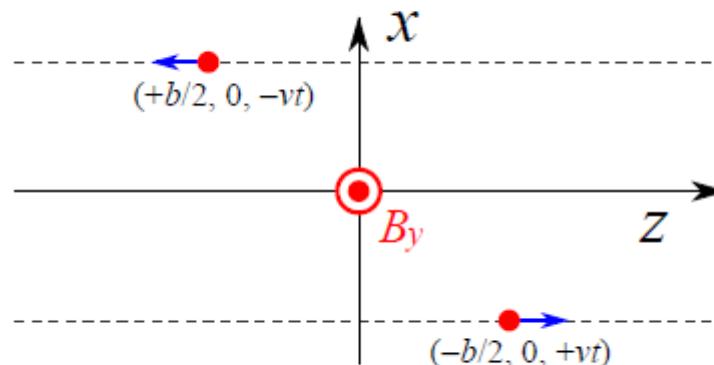


on top of the Quark-Gluon Plasma

Point-charge Approximation



Lienard-Wiechert potential



$$eB(t) = \frac{eB_0}{[1 + (t/t_0)^2]^{3/2}}$$

$$eB_0 = (47.6 \text{ MeV})^2 \left(\frac{1 \text{ fm}}{b} \right)^2 Z \sinh(Y), \quad t_0 = \frac{b}{2 \sinh(Y)}$$

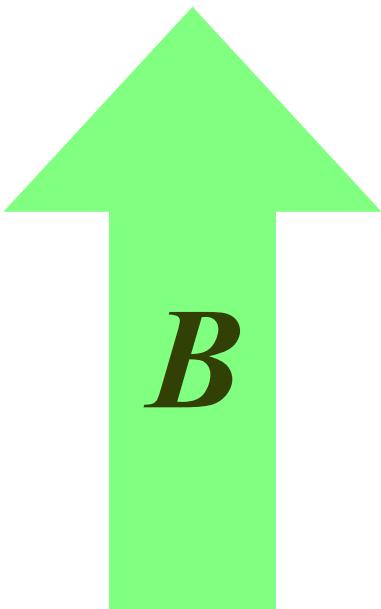
$\sim 10^{18}$ gauss

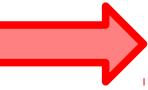
Interesting phenomena expected!

Discussed by Rafelski, Mueller, ... (~1976)

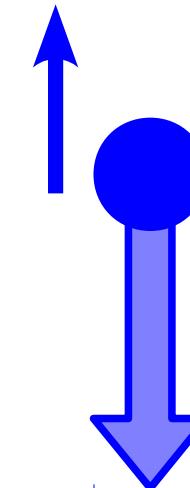
Chiral Magnetic Effect

Classical Picture



Right-handed Quark
= momentum 
parallel to
spin 

Left-handed Quark
= momentum 
anti-parallel to
spin 

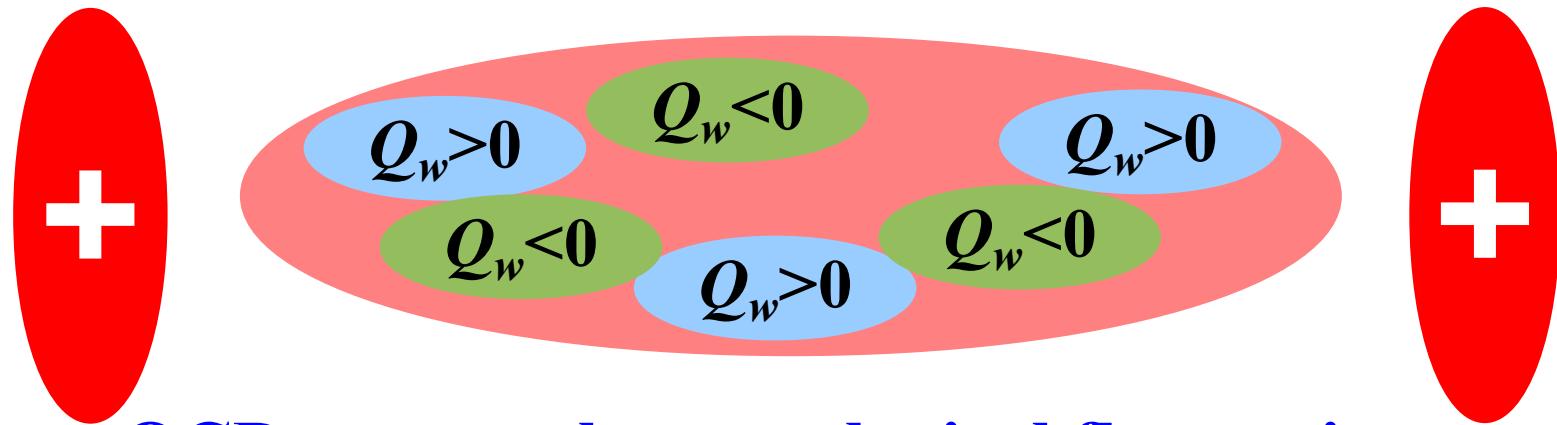


Finite Chirality
 $(\mathcal{P}$ breaking)

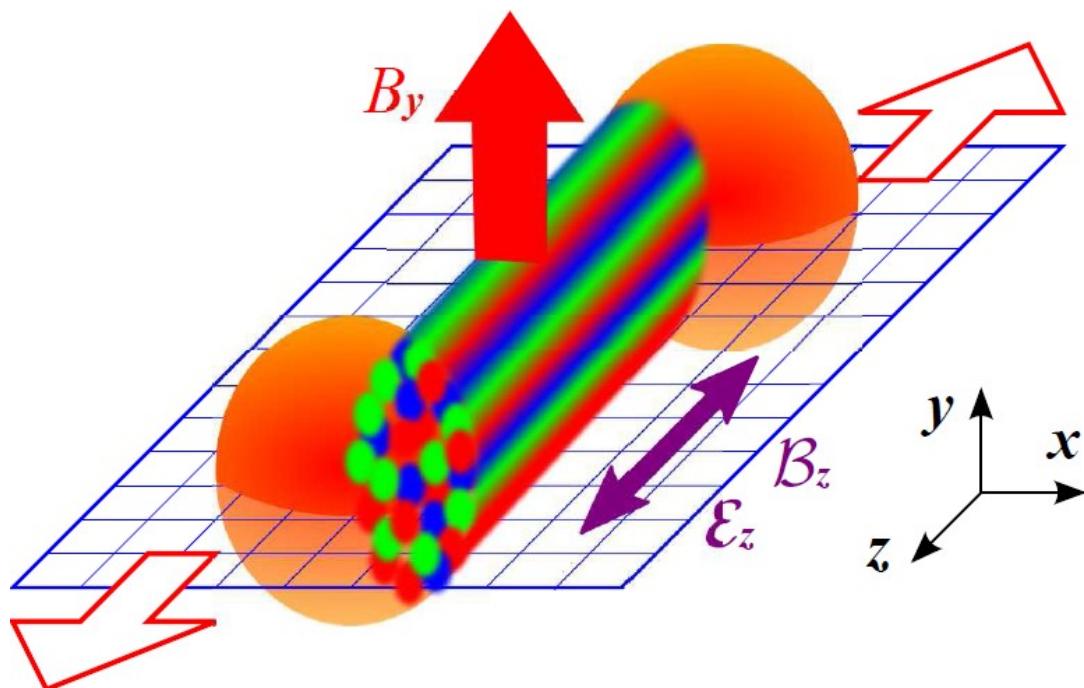

$$J \neq 0 \quad \text{if} \quad N_5 = N_R - N_L \neq 0$$

Kharzeev-McLerran-Warringa (2007)

Local Parity Violation (LPV)



QCD vacuum has topological fluctuations



**Initial condition of HIC
(Color Glass Condensate)
accommodates topological
fluctuations**

Current Generation through Anomaly



Chiral Magnetic Effect

$$j = N_c \sum_{\text{flavor}} \frac{\underline{q}_f^2 \mu_5}{2 \pi^2} B$$

Vilenkin (1980), Metlitski-Zhitnitsky, KF-Kharzeev-Warringa

Wess-Zumino-Witten Action



WZW term without U fields (contact term)

$$\begin{aligned}
 L_P = & \frac{N_c}{8 N_f \pi^2} \epsilon^{\mu\nu\rho\sigma} \left\{ \text{tr} \left[v_\mu \left(\partial_\nu v_\rho - \frac{i}{3} [v_\nu, v_\rho] \right) \right] \partial_\sigma \theta \right. \\
 & \left. + \text{tr} (a_\mu D_\nu a_\rho) \left[\frac{4}{3} \text{tr} (a_\sigma) + \partial_\sigma \theta \right] \right\} - \frac{N_c}{12 N_f^2 \pi^2} \text{tr} (a_\mu) \text{tr} (\partial_\nu a_\rho) \partial_\sigma \theta
 \end{aligned}$$

QED fields: $v_\mu = e Q$ $A_\mu = e \begin{pmatrix} 2/3 & 0 \\ 0 & -1/3 \end{pmatrix} A_\mu$

Kaiser-Leutwyler

Current Generation through WZW

$$\Gamma_p = \int \frac{N_c}{8 N_f \pi^2} \epsilon^{\mu\nu\sigma\rho} \text{tr} [\nu_\mu \partial_\nu \nu_\rho] \partial_\sigma \theta + \dots$$

$$j_z = \frac{d \Gamma_p}{d A_z} = \frac{N_c}{4 N_f \pi^2} \epsilon^{zxyt} \text{tr} (Q^2) B_z \partial_t \theta$$
$$= N_c \sum_f \frac{q_f^2 B_z}{2 \pi^2} \cdot \frac{\partial_t \theta}{2 N_f} \mu_5$$

produces non-zero chirality

Similar Effects

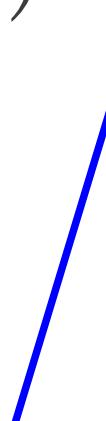
$$j_\mu \propto \epsilon_{\mu\nu\sigma\rho} (\partial^\nu \phi) F^\sigma \rho$$

Derivative of a pseudo-scalar quantity

η' condensate

pion condensates / profile

Strong θ angle



2nd-rank tensor
Field strength tensor
Angular momentum
Angular velocity

These effects under investigations in HIC

Summary



QCD phase diagram with chiral and deconfinement phase transitions is investigated:

- *High Temperature* – Phase transitions well understood from the zero-T properties of confinement.
- *High Baryon Density* – Inhomogeneous states favored near the phase boundary of homogeneous states.
- *Strong B Field* – Effects on the phase diagram not yet understood. Many interesting anomalous effects expected.

Experimental efforts focused on the baryon-rich matter and the visible effects of the strong B :

- Systematic fluctuation measurements to confirm the local parity violation / critical point / inhomogeneity